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Electronic structures and geometries of the XF$_3$ (X = Cl, Br, I, At) fluorides

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The potential energy surfaces of the group 17 XF$_3$ (X = Cl, Br, I, At) fluorides have been investigated for the first time with multiconfigurational wave function theory approaches. In agreement with experiment, bent T-shaped $C_{2v}$ structures are computed for ClF$_3$, BrF$_3$ and IF$_3$, while we predict that an average $D_{3h}$ structure would be experimentally observed for AtF$_3$. Electron correlation and scalar relativistic effects strongly reduce the energy difference between the $D_{3h}$ geometry and the $C_{2v}$ one, along the XF$_3$ series, and in the X = At case, spin-orbit coupling also slightly reduces this energy difference. AtF$_3$ is a borderline system where the $D_{3h}$ structure becomes a minimum, i.e. the pseudo Jahn-Teller effect is inhibited since electron correlation and scalar-relativistic effects create small energy barriers leading to the global $C_{2v}$ minima, although both types of effects interfere.

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I. INTRODUCTION

During the past decades, several computational works concerning the geometries of group 17 element XF\textsubscript{3} fluorides (where X = Cl, Br, I, At, or even element 117), were reported.\textsuperscript{1–9} These systems are not only of fundamental interest since they can be the subject of industrial applications. ClF\textsubscript{3} is a strong oxidizer, used as a fluorinating agent, and involved nowadays in the synthesis of uranium hexafluoride. BrF\textsubscript{3} is only used as a fluorinating agent in organic synthesis, while no application involving IF\textsubscript{3} is yet known. Several attempts to synthesize astatine fluorides turned out to be unsuccessful.\textsuperscript{10} The determination of the geometries of these known and hypothetical XF\textsubscript{3} systems is a task of fundamental interest in theoretical chemistry. The valence shell electron pair repulsion model\textsuperscript{11,12} supports bent T-shaped \( C_{2v} \) geometries (see Figure 1). Experimentally, such geometries were evidenced for ClF\textsubscript{3}\textsuperscript{13,14} and BrF\textsubscript{3}\textsuperscript{15} in the gas phase, and also in the condensed state of IF\textsubscript{3}.\textsuperscript{16} Regarding the (117)F\textsubscript{3} system, a planar \( D_{3h} \) structure is predicted from computations,\textsuperscript{4,5,9} while AtF\textsubscript{3} is a borderline system where the T-shaped \( C_{2v} \) and the \( D_{3h} \) structures are almost isoenergetic.\textsuperscript{2,4,5,9} For the XF\textsubscript{3} (X = Cl, Br, I) systems, such a \( D_{3h} \) structure is unstable due to the pseudo Jahn-Teller effect (PJTE).\textsuperscript{17} The detailed description of the PJTE in these cases was first given by Schwerdtfeger.\textsuperscript{2} In the AtF\textsubscript{3} case, the \( D_{3h} \) situation corresponds to a local minimum,\textsuperscript{5,9} meaning that the PJTE is in this case inhibited by other electronic factors.

![Figure 1](image)

FIG. 1. Ball and stick representations of T-shaped \( C_{2v} \) (left) and planar \( D_{3h} \) (right) XF\textsubscript{3} geometries. [Color code: maroon (X = Cl, Br, I, At), light blue (F)].

From a quantum mechanical perspective, non-relativistic (NR), scalar-relativistic (SR) and spin-orbit coupling (SOC) calculations have been performed\textsuperscript{1,2,4,5,9} in order to determine critical points on the ground state potential energy surfaces (PESs) of these XF\textsubscript{3} systems, as well as for determining energy differences between them, especially the one between the \( D_{3h} \) and the T-shaped \( C_{2v} \) structures, \( \Delta E^{D_{3h}-C_{2v}} \). However, for some systems, important
critical points of the PES are still not reported or discussed. Regarding most of the transition states (TSs), their nature (i.e. their connectivities with minima of the PESs) was generally not rigorously established. The most emblematic case is the AtF$_3$ system where important aspects are still a matter of debate in the literature. For instance, very recently Yang and Wang$^9$ stated that the T-shaped $C_{2v}$ structure is the only candidate to be observed through experiments. But one may think at first glance that their reasoning is a bit odd. Strong issues are related to the existence of the $D_{3h}$ structure as a true minimum, e.g. about the effects that inhibit the PJTE in this case. Different factors such as electron correlation, SR effects and SOC are able to significantly affect $\Delta E^{D_{3h}-C_{2v}}$. It is well known that numerous properties of At species are highly sensitive to SOC.$^{18-23}$ Thus, the most recent investigations$^{5,9}$ had focused on the relationship between SOC and the PJTE at the $D_{3h}$ structure. However, no definite statement has yet been made.$^9$ Indeed, one needs to consider the different factors not solely at the $D_{3h}$ structure, but also how they affect the PES in its close vicinity in order to reach a relevant conclusion. Furthermore, all previous theoretical studies made use of single-determinental methods while the reliability of such approaches has not been verified in difficult cases like AtF$_3$.

In this article, we aim at resolving the above-mentioned points through multiconfigurational wave function theory (WFT) calculations, via a contracted spin-orbit configuration interaction approach, and additional spin-orbit density functional theory (SO-DFT) calculations. Note that the two-step spin-orbit WFT approach have been already validated in iodine and astatine containing systems.$^{24,25}$ Since this work also deals with the ClF$_3$ and BrF$_3$ systems, we will start by validating the choice of the exchange-correlation functionals used to explore the PESs of interest. The obtained geometries of the critical points will be used for subsequent multiconfigurational WFT calculations, aiming at studying the dependence of the results on the degrees of freedom associated to this methodology. Finally, based on multiconfigurational calculations, we will discuss the role of electron correlation, scalar relativistic effects, and SOC along the XF$_3$ series in order to clarify and complement the conclusions of previous investigations$^{4,5,9}$ on the respective role of these effects on the PESs of the studied systems.
II. METHODOLOGY

A. DFT calculations

The critical points of the PESs corresponding to the ground state of the XF₃ (X = Cl–At) systems were optimized using NR-DFT (X = Cl), SR-DFT and SO-DFT (X = Br–At) approaches. NR- and SR-DFT calculations have been performed using the GAUSSIAN09 program package, while all the SO-DFT ones were carried out with the NWChem program package. Vibrational harmonic frequencies have been determined to establish the nature of the different critical points. Afterwards, the connectivities between TSs and minima were established through the computation of intrinsic reaction coordinate (IRC) paths at NR or SR levels.

The M06-2X and PBE0 hybrid exchange-correlation functionals were considered, since these two functionals showed good performance on the XF₃ (X = I–117) fluorides. In order to further discuss the role of the fraction of “exact” exchange on the performance of DFT functionals on these systems, this fraction has been also varied in a hybrid PBE functional, PBEₓ(x), where x denotes the percentage of exact exchange.

In the SR-DFT and SO-DFT calculations, energy-consistent small-core ECPₙMDF pseudopotentials were used to mimic the effects of the n = 60 core electrons of At, n = 28 core electrons of I, and n = 10 core electrons of Br. Note that in the SO-DFT calculations, the pseudopotentials account for both scalar relativistic effects and SOC. For describing the explicitly treated electrons (25 in the neutral atoms), the aug-cc-pVTZ-PP-2c basis sets were used for the At and I atoms, while the aug-cc-pVTZ-PP one was used for Br. All-electron aug-cc-pVTZ basis sets were used for the Cl and F atoms.

B. WFT calculations

All the WFT calculations were performed with the ORCA program package. Regarding the ClF₃ and BrF₃ systems, reference values for ΔE⁢D₃h−C₂v were obtained with the CCSD(T) method by means of geometry optimizations. For ClF₃, the relativistic effects were neglected and TZVP basis sets were used. For BrF₃, scalar relativistic effects were taken into account via the use of the Douglas-Kroll-Hess Hamiltonian and accordingly, the segmented all-electron relativistically contracted SARC-TZVP-DKH basis set were used.
Two-step WFT calculations were then performed. In the first step, a set of spin-orbit free (SOF) states is computed at a multiconfigurational level, here using the state-averaged complete active space self-consistent field (SA-CASSCF) approach.\textsuperscript{42,43} The calculations were performed on the previously optimized (SO-)DFT structures. In the second step, the spin components of the SOF states are used to build a state-interaction matrix, consisting of correlated energies on the diagonal, and off-diagonal SOC matrix elements computed within a mean-field approximation.\textsuperscript{44,45} An important degree of freedom concerns the method used to compute the correlated energies, which can be SA-CASSCF or the \( n \)-electron valence state second-order perturbation theory (NEVPT2) method.\textsuperscript{46} The latter accounts for more dynamic correlation and usually does not suffer from the occurrence of intruder states. The SARC-TZVP-DKH basis sets\textsuperscript{41,47} were used for all atoms in these calculations, and scalar relativistic effects were accounted for via the Douglas-Kroll-Hess Hamiltonian.\textsuperscript{38–40}

Regarding the CASSCF calculations, two active spaces have been considered: (i) a minimal one (in view of SOC calculations) comprising only 2 electrons and the \( np \) orbitals of the central atom, leading to CASSCF(2/3) calculations, and (ii) a larger one comprising 20 electrons and the \( np \) and \( 2p \) orbitals of the central and F atoms, leading to CASSCF(20/12) calculations. In the SA-CASSCF(2/3) calculations, two state-averaging schemes have been considered: the full one consisting of the \( M_{S_{\text{max}}} \) components of the 6 singlet (S) and 3 triplet (T) spin states (6S3T), and a subset of it, containing the ground SOF state plus all the essentially singly-excited states, \( i.e. \) 3S and 2T states (3S2T). For the SA-CASSCF(20/12) calculations, only one averaging scheme has been considered, namely 21S20T. Note that this averaging-scheme fulfills the following conditions for all the studied XF\(_3\) systems, (i) it does not artificially break the symmetry in the \( D_{3h} \) structures, (ii) it includes all the essentially singly excited states, and (iii) the enlargement of the state-averaging space and of the corresponding state-interaction space hardly affect the computed SOC ground state energy stabilizations.

Furthermore, specific Hartree-Fock (HF) and NEVPT2 calculations were performed at NR and SR levels in order to highlight the role of scalar relativistic effects, electron correlation and SOC on the stability of some structures. These calculations made use of fully uncontracted SARC-TZVP-DKH basis sets in order to avoid any bias introduced by the contraction scheme.
III. RESULTS AND DISCUSSION

A. DFT study

We start our analysis by looking at the performance of the M06-2X and PBE0 functionals in determining the geometries of the $D_{3h}$ and T-shaped $C_{2v}$ structures, as well as the energy difference between them ($\Delta E^{D_{3h}-C_{2v}}$). With both functionals, the T-shaped $C_{2v}$ structures are minima of the PESs for all the studied XF$_3$ systems. The $D_{3h}$ structure is a second-order saddle point in the X = Cl–I cases, with one doubly-degenerate $e'$ imaginary frequency. In contrast, the $D_{3h}$ structure is found to be a (weak) local minimum in the AtF$_3$ case. These results are compared to experimental data, when available, as well as to (SO-)CCSD(T) reference calculations (see Table I). Since the CCSD(T) equilibrium geometries are in very good agreement with the experimental ones for the ClF$_3$ and BrF$_3$ T-shaped $C_{2v}$ structures, both sets of results can be alternatively used to assess the quality of the (SO-)DFT results. Table I reveals that both the PBE0 and M06-2X functionals lead to geometries in very good agreement with the (SO-)CCSD(T) ones, and when available, with the experimental data. For instance, the obtained deviation with the PBE0 functional regarding the equilibrium bond lengths with respect to the (SO-)CCSD(T) ones is less than 2% for each studied system. However, important discrepancies with the CCSD(T) results appear in some cases for $\Delta E^{D_{3h}-C_{2v}}$. While (SO-)PBE0 gives reasonable energy differences compared to (SO-)CCSD(T) for the whole XF$_3$ series, the (SO-)M06-2X results deviate more and more from X = At to X = Cl. For ClF$_3$, the $\Delta E^{D_{3h}-C_{2v}}$ value computed with the M06-2X functional is too large by $\sim$30 kJ mol$^{-1}$ with respect to the CCSD(T) value. One can conclude that the PBE0 functional is overall more adequate than the M06-2X one for studying PESs in the XF$_3$ (X = Cl–At) series.

In order to highlight the role of the fraction $x$ of exact exchange in hybrid functionals, we monitored the $r_e$ bond length ($D_{3h}$ case) and $\Delta E^{D_{3h}-C_{2v}}$ while varying $x$ in PBEh($x$) functional, from 20 to 100 % (see Figure 2). Using B3LYP as the reference functional, Yang and Wang$^9$ had performed a similar study but restricted to the case of AtF$_3$. We find similar trends in the whole XF$_3$ series, i.e. both $r_e$ and $\Delta E^{D_{3h}-C_{2v}}$ diminish with the fraction of exact exchange. When comparing to the (SO-)CCSD(T) results, we find, for instance in the case of AtF$_3$, a different optimal fraction, $\sim$35 % for $r_e$ and $\sim$40 % for $\Delta E^{D_{3h}-C_{2v}}$, with
TABLE I. Geometrical parameters of the $D_{3h}$ and T-shaped $C_{2v}$ structures computed at various levels of theory, and associated energy differences.

<table>
<thead>
<tr>
<th>System</th>
<th>Method</th>
<th>$D_{3h}$</th>
<th>T-shaped $C_{2v}$</th>
<th>$\Delta E^{D_{3h}-C_{2v}}$ (kJ mol$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>$r_e$(Å)</td>
<td>$r_{eq}$(Å)</td>
<td>$r_{ax}$(Å)</td>
<td>$\alpha$(°)</td>
</tr>
<tr>
<td>ClF$_3$</td>
<td>NR-M06-2X</td>
<td>1.713</td>
<td>1.592</td>
<td>1.696</td>
</tr>
<tr>
<td></td>
<td>NR-PBE0</td>
<td>1.728</td>
<td>1.603</td>
<td>1.706</td>
</tr>
<tr>
<td></td>
<td>SR-CCSD(T)</td>
<td>1.738</td>
<td>1.605</td>
<td>1.703</td>
</tr>
<tr>
<td></td>
<td>Expt.$^{13}$</td>
<td>-</td>
<td>1.598</td>
<td>1.698</td>
</tr>
<tr>
<td>BrF$_3$</td>
<td>SO-M06-2X</td>
<td>1.815</td>
<td>1.715</td>
<td>1.802</td>
</tr>
<tr>
<td></td>
<td>SO-PBE0</td>
<td>1.825</td>
<td>1.725</td>
<td>1.812</td>
</tr>
<tr>
<td></td>
<td>SR-CCSD(T)</td>
<td>1.828</td>
<td>1.725</td>
<td>1.811</td>
</tr>
<tr>
<td></td>
<td>Expt.$^{15}$</td>
<td>-</td>
<td>1.721</td>
<td>1.810</td>
</tr>
<tr>
<td>IF$_3$</td>
<td>SO-M06-2X</td>
<td>1.956</td>
<td>1.866</td>
<td>1.940</td>
</tr>
<tr>
<td></td>
<td>SO-PBE0</td>
<td>1.965</td>
<td>1.878</td>
<td>1.950</td>
</tr>
<tr>
<td></td>
<td>SO-CCSD(T)$^9$</td>
<td>1.941</td>
<td>1.853</td>
<td>1.926</td>
</tr>
<tr>
<td>AtF$_3$</td>
<td>SO-M06-2X</td>
<td>2.059</td>
<td>1.993</td>
<td>2.059</td>
</tr>
<tr>
<td></td>
<td>SO-PBE0</td>
<td>2.068</td>
<td>2.012</td>
<td>2.068</td>
</tr>
<tr>
<td></td>
<td>SO-CCSD(T)$^9$</td>
<td>2.046</td>
<td>1.987</td>
<td>2.046</td>
</tr>
</tbody>
</table>

respect to the $\sim 55$% obtained in Ref. 9. This discrepancy is due to the use of another reference functional. Therefore, the good performance of hybrid functionals for computing geometries and energy differences in the XF$_3$ series is not only related to the fraction of exact exchange since different functionals with very different fractions of exact exchange can lead to a similar accuracy. From Figure 2, one would conclude that the recommended fractions of exact exchange would range from 20–30% (ClF$_3$ case) to 35–45% (IF$_3$ case) for PBEh($x$). Overall, one could recommend to use the PBEh(35) functional. However, not much improvement would be reached compared to the standard PBE0 functional ($x=25$), which has therefore been used in the rest of this work.

We now question the nature of the Y-shaped $C_{2v}$ TSs on the PESs of the XF$_3$ systems, that were previously reported by Yang and Wang.$^9$ Table II displays the computed geometrical parameters using both functionals, as well as the energy difference of these Y-shaped TSs.
FIG. 2. Evolution of the energy differences between the $C_{2v}$ and $D_{3h}$ structures ($\Delta E^{D_{3h} - C_{2v}}$) and of the $D_{3h}$ bond lengths ($r_e$) as a function of the fraction of exact exchange ($x$) introduced in the PBEh($x$) functional.

with respect to the T-shaped $C_{2v}$ global minima. The comparison with the $D_{3h}$ structure shows for the ClF$_3$, BrF$_3$ and IF$_3$ cases that the $r_{ax}$ bond lengths are smaller than $r_e$, the $r_{eq}$ bond lengths are larger than $r_e$ and the $\alpha$ bond angles are larger than 120°. Such Y-shaped TSs are expected to be involved in the pseudo-rotations between two equivalent T-shaped $C_{2v}$ global minima, as suggested in the ClF$_3$\textsuperscript{1} and IF$_3$\textsuperscript{9} cases. We confirm this hypothesis through IRC calculations as displayed in Figure 3. Note that going from the $D_{3h}$ to the T-shaped $C_{2v}$ structure via a Y-shaped $C_{2v}$ TS does not involve a drastic change in the electronic configuration.

Regarding the AtF$_3$ system, a Y-shaped $C_{2v}$ TS also exists but its structure exhibits an $\alpha$ bond angle smaller than 120°, an $r_{ax}$ bond length larger than $r_e$, and an $r_{eq}$ bond length
FIG. 3. Energies along the IRC paths starting from Y-shaped $C_{2v}$ TSs, computed at NR- or SR-PBE0 levels (the T-shaped $C_{2v}$ energies are chosen as the zeroes of energy).

TABLE II. Geometrical parameters of the Y-shaped $C_{2v}$ TSs computed at various levels of theory, and associated energy differences $\Delta E^{TS-C_{2v}}$. 

<table>
<thead>
<tr>
<th>System</th>
<th>Method</th>
<th>$r_{eq}$ (Å)</th>
<th>$r_{ax}$ (Å)</th>
<th>$\alpha$ (°)</th>
<th>$\Delta E^{TS-C_{2v}}$ (kJ mol$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>ClF$_3$</td>
<td>NR-M06-2X</td>
<td>1.770</td>
<td>1.676</td>
<td>130.9</td>
<td>121.5</td>
</tr>
<tr>
<td></td>
<td>NR-PBE0</td>
<td>1.762</td>
<td>1.703</td>
<td>128.7</td>
<td>89.3</td>
</tr>
<tr>
<td>BrF$_3$</td>
<td>SO-M06-2X</td>
<td>1.848</td>
<td>1.793</td>
<td>128.7</td>
<td>90.8</td>
</tr>
<tr>
<td></td>
<td>SO-PBE0</td>
<td>1.843</td>
<td>1.814</td>
<td>125.6</td>
<td>59.9</td>
</tr>
<tr>
<td>IF$_3$</td>
<td>SO-M06-2X</td>
<td>1.981</td>
<td>1.940</td>
<td>127.8</td>
<td>77.9</td>
</tr>
<tr>
<td></td>
<td>SO-PBE0</td>
<td>1.978</td>
<td>1.958</td>
<td>124.4</td>
<td>54.9</td>
</tr>
<tr>
<td></td>
<td>SO-CCSD(T)$^9$</td>
<td>1.960</td>
<td>1.930</td>
<td>126.4</td>
<td>67.3</td>
</tr>
<tr>
<td>AtF$_3$</td>
<td>SO-M06-2X</td>
<td>2.053</td>
<td>2.060</td>
<td>117.8</td>
<td>11.1</td>
</tr>
<tr>
<td></td>
<td>SO-PBE0</td>
<td>2.047</td>
<td>2.075</td>
<td>108.7</td>
<td>3.1</td>
</tr>
<tr>
<td></td>
<td>SO-CCSD(T)$^9$</td>
<td>2.039</td>
<td>2.049</td>
<td>116.0</td>
<td>7.8</td>
</tr>
</tbody>
</table>

smaller than $r_e$. Furthermore, this TS is (slightly) higher in energy than the $D_{3h}$ structure, in contrast to what is observed for ClF$_3$, BrF$_3$ and IF$_3$. Hence, this TS exhibits a different nature and it was suggested to connect T-shaped $C_{2v}$ and $D_{3h}$ minima by Yang and Wang.$^9$ Note that differences appear between the results presented in Table II and the energy profiles displayed in Figure 3 since SOC is neglected in these IRC calculations. In the AtF$_3$ system,
FIG. 4. $C_s$ TS structure of AtF$_3$ with displacement vectors traced for the imaginary frequency (left) and energy along the IRC path starting from the $C_s$ TS and computed at SR-PBE0 level (right). The energy of the T-shaped $C_{2v}$ minimum is chosen as the zero of energy.

SOC notably stabilizes the Y-shaped $C_{2v}$ and the $D_{3h}$ structures compared to the T-shaped one.

In the literature, the potential existence of TSs connecting T-shaped $C_{2v}$ minima in AtF$_3$ was not considered. At SR levels, we have identified a $C_s$ TS with an out-of-plane distorted geometry (see Figure 4). This $C_s$ TS actually connects two T-shaped $C_{2v}$ minima, as shown by the IRC profile (see Figure 4). SOC significantly stabilizes the TS and modifies its geometrical parameters close to a $C_{3v}$ structure at the SO-M06-2X level, and leads to the practical impossibility of converging on a $C_s$ structure at the SO-PBE0 level. The optimized SR-DFT and SO-DFT geometrical parameters of the $C_s$ TS are compiled in Table III.

In close connection with the $C_s$ TS structure, a second-order saddle point belonging to the $C_{3v}$ symmetry point group was evidenced. The doubly-degenerate imaginary frequencies are 57i and 67i cm$^{-1}$ at the SO-M06-2X and SO-PBE0 levels, respectively. Note that in the limit of a doubly-degenerate zero-frequency normal mode, this $C_{3v}$ structure would correspond to a “monkey” saddle point associated with the trifurcation toward the three possible $C_s$ TS structures. Such situation was for instance suggested in (tetracoordinated) AX$_4$ systems, even if the existence of “monkey” saddle points can be seriously questioned since
it would violate some rules. The fact that we failed to optimize the $C_s$ TS structure at the SO-PBE0 level could be related to the near degeneracy with the $C_{3v}$ structure. The near degeneracy between the $C_s$ and $C_{3v}$ geometries is also an indication for the existence of an effective “monkey” $C_{3v}$ saddle point in AtF$_3$, even if no definitive conclusion can actually be given on this point.

**TABLE III.** AtF$_3$ geometrical parameters of the $C_s$ TS computed at SO-DFT and SR-DFT (in parenthesis) levels of theory, corresponding imaginary harmonic frequency, and associated energy difference $\Delta E^{TS-C_{2v}}$.

<table>
<thead>
<tr>
<th>Method</th>
<th>$r_{in}$ (Å)</th>
<th>$r_{out}$ (Å)</th>
<th>$\nu_1$ (cm$^{-1}$)</th>
<th>$\Delta E^{TS-C_{2v}}$ (kJ mol$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>SO-M06-2X</td>
<td>2.061 (1.957)</td>
<td>2.062 (2.063)</td>
<td>106i (314i)</td>
<td>156.4 (256.0)</td>
</tr>
<tr>
<td>SO-PBE0</td>
<td>(1.969)</td>
<td>(2.078)</td>
<td>(308i)</td>
<td>(242.8)</td>
</tr>
</tbody>
</table>

Schematic representations of the connectivity maps for the ground state PESs of the studied XF$_3$ systems are shown in Figure 5. For X = Cl–I, the $D_{3h}$ structure is a second-order saddle point and three equivalent T-shaped $C_{2v}$ global minima are connected via Y-shaped $C_{2v}$ TSs. For AtF$_3$, the $D_{3h}$ structure is a “weak” local minimum, while three equivalent T-shaped $C_{2v}$ structures correspond to global minima of the PES. Between the $D_{3h}$ and each T-shaped $C_{2v}$ structure, there is a very small energy barrier associated to a Y-shaped $C_{2v}$ TS. At SR levels, each pair of T-shaped $C_{2v}$ minima are connected via a $C_s$ TS leading to a large energy barrier. When SOC is included in M06-2X calculations, the $C_s$ structure is brought very close in terms of geometry and energy to a more symmetric $C_{3v}$ one. Note that dissociation energies corresponding to the formation of either XF + F$_2$ or XF$_2$ + F are reported in Table SI. All the critical points reported in Figure 5 are found lower in energy than the lowest dissociation limit for each XF$_3$ system.

Finally, we find that the inclusion of entropic and thermal contributions significantly reduces the difference in stability between the $D_{3h}$ and $C_{2v}$ minima. At $T = 298$ K and $P = 1$ atm, the difference in free energy is 3.6 and 0.3 kJ mol$^{-1}$ at the SO-M06-2X and SO-PBE0 levels, respectively. Given the accuracy of the used approaches, it becomes impossible to conclude on the nature of the AtF$_3$ most stable structure ($C_{2v}$ or $D_{3h}$). However, since the free energy difference between these structures and the associated energy barriers are so small, it would be impossible to experimentally distinguish between these two
structures, and therefore, an average $D_{3h}$ geometry would be seen in the experiments, if they were possible.

![Diagram of connectivity maps between critical points](image)

**FIG. 5.** Schemes of the connectivity maps between the important critical points of the PESs of the XF$_3$ systems, X = Cl–I (left), X = At (right). [Color code: blue is used for minima, green for TSs and red for second-order saddle points.]

### B. WFT study

The dependence of the results with respect to computational degrees of freedom was first assessed by means of single-point calculations on previously obtained DFT geometries. Since in two-step WFT calculations the SOC is introduced *a posteriori*, it is straightforward to distinguish SOF and SOC contributions to $\Delta E^{D_{3h}-C_{2v}}$. In Tables SII–SIV and IV, SOC stabilizations of the T-shaped $C_{2v}$ and the $D_{3h}$ structures, namely $E^{C_{2v}}_{SOC}$ and $E^{D_{3h}}_{SOC}$ are reported. It is therefore possible to quantify SOC effects ($\Delta E^{D_{3h}-C_{2v}}_{SOC}$) on $\Delta E^{D_{3h}-C_{2v}}$.

Tables SII and SIII collect energy parameters computed at the NEVPT2(2/3) and NEVPT2(2/3)+SOC levels of theory, with 6S3T and 3S2T averaging schemes. As expected, both schemes lead to quite similar results. At these levels, $\Delta E^{D_{3h}-C_{2v}}$ is rather constant in the series from $X = \text{Cl}$ to $X = \text{I}$, while it is much smaller in the AtF$_3$ case. According to the reference CCSD(T) results, a significant decrease of $\Delta E^{D_{3h}-C_{2v}}$ is expected from ClF$_3$ to BrF$_3$, which is not reproduced at the considered NEVPT2(2/3) levels. This fact is probably due to the size of the active space, which may not be large enough to well describe covalency...
in the states of interest. As expected, the absolute values of $E_{SOC}^{C_{2v}}$ and $E_{SOC}^{D_{3h}}$ increase with the atomic number. The SOC stabilization of $C_{2v}$ and $D_{3h}$ structures being practically equivalent from $X = \text{Cl}–\text{I}$, no significant SOC contribution to $\Delta E_{SOC}^{D_{3h}–C_{2v}}$ is observed in Tables SII and SIII. The AtF$_3$ case is more ambiguous since $\Delta E_{SOC}^{D_{3h}–C_{2v}}$ is of the same order of magnitude as the NEVPT2(2/3) $\Delta E_{SOC}^{D_{3h}–C_{2v}}$, leading to an almost vanishing $\Delta E_{SOC}^{D_{3h}–C_{2v}}$ value at the NEVPT2(2/3)+SOC level. It is thus clear that increasing the active space is actually needed for determining which of the $C_{2v}$ and the $D_{3h}$ structures is lower in energy at the NEVPT2+SOC level of theory.

SA-CASSCF(20/12) and NEVPT2(20/12) results obtained with active spaces which also include the $2p$ orbitals (and electrons) of fluorine atoms are presented in Tables SIV and IV. The comparison between the two sets of data highlights the importance of the electron correlation that was not treated within the active space. While this correlation enlarges $\Delta E_{SOC}^{D_{3h}–C_{2v}}$ in the ClF$_3$ and BrF$_3$ cases, it markedly diminishes it in the IF$_3$ and AtF$_3$ cases. At the SA-CASSCF(20/12)+SOC level, the expected trend regarding $\Delta E_{SOC}^{D_{3h}–C_{2v}}$ values within the series is not reproduced, while it appears correct at the NEVPT2(20/12)+SOC level, with values in good agreement with the (SO-)CCSD(T) ones. The SOC contribution to the energy lowering of the $C_{2v}$ and $D_{3h}$ structures is negligible for ClF$_3$ and BrF$_3$, is moderate for IF$_3$ ($\sim 5$ kJ mol$^{-1}$), and appears quite important for AtF$_3$ ($\sim 60$ kJ mol$^{-1}$). As previously noted, SOC does not contribute much to $\Delta E_{SOC}^{D_{3h}–C_{2v}}$, since this coupling stabilizes in a very similar way the $C_{2v}$ and $D_{3h}$ minima. However, the $\Delta E_{SOC}^{D_{3h}–C_{2v}}$ value in the AtF$_3$ case, which is about $–0.7$ kJ mol$^{-1}$, appears somewhat small compared to previous DFT calculations: the energy difference between the $D_{3h}$ and T-shaped $C_{2v}$ minima is reduced by 8.8(11.6) kJ mol$^{-1}$ when comparing SO-PBE0(SO-M06-2X) results. This issue is most probably related to the neglect of the geometry relaxation when computing the NEVPT2(20/12) and NEVPT2(20/12)+SOC energies (SO-PBE0 geometries are used).

It is worth noting for the whole series that in the SA-CASSCF(20/12) ground-state wave functions, the weights of the HF configurations are always larger than 80%. Due to the use of state-averaged orbitals, the SOF wave functions may look more multiconfigurational than what they should actually be. This means that these ground SOF states are essentially single-configurational in all the series. In addition, regarding the SOC wave functions, they are dominated by the ground-SOF-singlet spin component. In the AtF$_3$ case, where SOC effects are the largest ones, the $|M_S| = 1$ components of the first two excited SOF triplet
states correspond to 16% of the ground SOC wave function for the \(D_{3h}\) minimum, while the ground-SOF-singlet spin component weights 83%. We definitively conclude that the ground SOC wave functions are essentially single-configurational in this series, which justifies the use of single-determinantal or single-reference approaches such as the SO-DFT and SO-CCSD(T) ones.

**TABLE IV.** Contributions (in kJ mol\(^{-1}\)) to \(\Delta E_{D_{3h}-C_{2v}}\) computed at the NEVPT2 and NEVPT2+SOC levels using SA-CASSCF(20/12) references and 21S20T state-averaging and state-interaction spaces. The reference (SO-)CCSD(T) \(\Delta E_{D_{3h}-C_{2v}}\) values are also recalled for convenience.

<table>
<thead>
<tr>
<th>System</th>
<th>NEVPT2</th>
<th>NEVPT2+SOC</th>
<th>(SO-)CCSD(T)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>(\Delta E_{D_{3h}-C_{2v}})</td>
<td>(E_{D_{3h}}^{\text{SOC}})</td>
<td>(E_{C_{2v}}^{\text{SOC}})</td>
</tr>
<tr>
<td>ClF(_3)</td>
<td>88.6</td>
<td>-0.1</td>
<td>-0.1</td>
</tr>
<tr>
<td>BrF(_3)</td>
<td>72.0</td>
<td>-1.2</td>
<td>-1.0</td>
</tr>
<tr>
<td>IF(_3)</td>
<td>66.2</td>
<td>-6.1</td>
<td>-5.4</td>
</tr>
<tr>
<td>AtF(_3)</td>
<td>7.4</td>
<td>-58.8</td>
<td>-58.1</td>
</tr>
</tbody>
</table>

**FIG. 6.** Scheme representing the two different pathways used to define the contributions of scalar relativistic effects (SRE) and electron correlation (EC) on \(\Delta E_{D_{3h}-C_{2v}}\).

In order to shed light on important factors affecting the \(D_{3h} - C_{2v}\) energy difference within
the series, namely scalar relativistic effects (SRE) and electron correlation (EC), NR-HF, SR-HF, NR-NEVPT2(20/12) and SR-NEVPT2(20/12) calculations have been performed with uncontracted SARC-TZVP-DKH basis sets (see Table V). It allows us to define SRE and EC contributions to the SR $\Delta E_{D_{3h}-C_{2v}}$ values through two paths: (i) Path\textsubscript{1} in which SRE are first introduced at the non-correlated HF level, prior to introducing EC at the SR level, and (ii) Path\textsubscript{2} in which EC is introduced in the absence of SRE and SRE are later introduced at the NEVPT2 correlated level (see Figure 6). If SRE and EC were additive, both paths would yield the same SRE and EC contributions to $\Delta E_{D_{3h}-C_{2v}}$. One concludes that it is not the case by looking at the data reported in Table V, \textit{i.e.} SRE and EC interfere. In the X = Cl–I cases, $\Delta E_{D_{3h}-C_{2v}}$ is more affected by EC than by SRE. In contrast, an opposite behavior is observed in the AtF\textsubscript{3} case, which furthermore displays a particularly large interference between SRE and EC: the sum of SRE(HF) and EC(NR) contributions is about $\sim 20$ kJ mol\textsuperscript{−1} over the sum of the EC(SR) and SRE(NEVPT2) contributions. Note that although the NR-HF $\Delta E_{D_{3h}-C_{2v}}$ values are too large compared to the SR-NEVPT2 ones, they already exhibit a correct trend within the series. Indeed, the simultaneous treatment of both SRE and EC affects in a similar way the $\Delta E_{D_{3h}-C_{2v}}$ values (with respect to the NR-HF ones), with only a small noticeable difference for the ClF\textsubscript{3} case (SRE and EC decrease more $\Delta E_{D_{3h}-C_{2v}}$ than in the other cases).

### Table V

<table>
<thead>
<tr>
<th>XF\textsubscript{3}</th>
<th>NR-HF $\Delta E_{D_{3h}-C_{2v}}$</th>
<th>Path\textsubscript{1}</th>
<th>Path\textsubscript{2}</th>
<th>SR-NEVPT2 $\Delta E_{D_{3h}-C_{2v}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>SRE(HF)</td>
<td>EC(SR)</td>
<td>EC(NR)</td>
<td>SRE(NEVPT2)</td>
</tr>
<tr>
<td>ClF\textsubscript{3}</td>
<td>185.8</td>
<td>−92.6</td>
<td>−93.6</td>
<td>−3.1</td>
</tr>
<tr>
<td>BrF\textsubscript{3}</td>
<td>149.7</td>
<td>−63.1</td>
<td>−65.9</td>
<td>−11.9</td>
</tr>
<tr>
<td>IF\textsubscript{3}</td>
<td>142.5</td>
<td>−46.3</td>
<td>−51.1</td>
<td>−24.5</td>
</tr>
<tr>
<td>AtF\textsubscript{3}</td>
<td>80.5</td>
<td>−22.1</td>
<td>−31.5</td>
<td>−42.9</td>
</tr>
</tbody>
</table>

In order to assess the existence of the low TS connecting $D_{3h}$ and $C_{2v}$ minima in the AtF\textsubscript{3} system, one has to go beyond single-point calculations. We have considered scans of the ground state PES between the $D_{3h}$ and $C_{2v}$ structures \textit{via} NEVPT2(20/12) and NEVPT2(20/12)+SOC calculations. Three variables naturally arise, namely $\alpha$, $r_{ax}$ and $r_{eq}$.
FIG. 7. Representation of the NEVPT2(20/12) (left) and NEVPT2(20/12)+SOC (right) ground state PESs of AtF$_3$. The energy scale is expressed in kJ mol$^{-1}$ and the T-shaped $C_{2v}$ energies are chosen as the zeroes of energy.

which were all varied in contrast to previous studies.\cite{2,5,9} In order to approximately determine the minimum energy path, a non-uniform grid has been defined by interpolations between the $D_{3h}$ and $C_{2v}$ structures. In Figure 7, contour plots of the resulting surfaces are represented. Note that for each $r_{eq}$ value corresponds a different $r_{ax}$ value. The approximate structures of the $D_{3h}$ and $C_{2v}$ minima at NEVPT2+SOC level do not differ much from the ones optimized at the SO-PBE0 level. As can also be seen, the Y-shaped $C_{2v}$ TS is evidenced between the $D_{3h}$ and $C_{2v}$ minima at both NEVPT2 and NEVPT2+SOC levels. Its approximate structure complies with the previous findings: $\alpha$ smaller than 120°, $r_{ax}$ larger than $r_e$ and $r_{eq}$ smaller than $r_e$ (not explicitly shown on Figure 7). Notable SOC effects on the PES are evidenced by the comparison of the NEVPT2 and NEVPT2+SOC scans displayed in Figure 7. It is clear that the energy difference between the $D_{3h}$ and $C_{2v}$ minima, and to a lesser extend, the one between the Y-shaped $C_{2v}$ TS and the $C_{2v}$ minimum are decreased, i.e. the PES is “flattened” by SOC. This does agree with the reference SO-CCSD(T) calculations, where SOC reduces about 8 kJ mol$^{-1}$ the energy difference between the $D_{3h}$ and $C_{2v}$ minima.\cite{9}

Finally, we have analyzed the factors affecting the energy difference between the Y-shaped $C_{2v}$ TS and the $D_{3h}$ minimum by means of single-point calculations, performed on the SO-PBE0 geometries. The results gathered in Table VI show that SRE and EC are the main contributions and are both responsible for the existence of the Y-shaped TS. Indeed, SRE and
EC tend to enlarge the energy difference, resulting in a $C_{2v}$ TS lying at higher energy than the $D_{3h}$ minimum. As for the $\Delta E^{D_{3h}-C_{2v}}$ energy difference, some interference appear between SRE and EC which actually lead to a smaller energy difference than if the effects were additive. SOC does not seem to affect much the energy difference, an energy increase of only 0.3 kJ mol$^{-1}$ is calculated. The latter prediction is quite consistent with previous DFT results: the comparison of SR-PBE0 (SR-M06-2X) and SO-PBE0 (SO-M06-2X) results, points out that SOC increases the energy difference between the Y-shaped $C_{2v}$ TS and the $D_{3h}$ minimum by 0.6 (0.1) kJ mol$^{-1}$. It is worth noting that previous SO-CCSD(T) calculations led to a small decrease of about 1 kJ mol$^{-1}$. Thus, it comes out that spin-dependent relativistic effects are definitely small and somewhat bland on the relative energy of the Y-shaped $C_{2v}$ TS with respect to the one of the $D_{3h}$ minimum.

### TABLE VI. Energy difference (in kJ mol$^{-1}$) between the Y-shaped TS and the $D_{3h}$ minimum of AtF$_3$ ($\Delta E^{TS-D_{3h}}$) computed at different levels of theory (SO-PBE0 geometries are considered).

<table>
<thead>
<tr>
<th>XF$_3$</th>
<th>NR-HF</th>
<th>SR-HF</th>
<th>NR-NEVPT2</th>
<th>SR-NEVPT2</th>
<th>SR-NEVPT2+SOC</th>
</tr>
</thead>
<tbody>
<tr>
<td>AtF$_3$</td>
<td>$-12.0$</td>
<td>$-2.5$</td>
<td>$-6.2$</td>
<td>$1.8$</td>
<td>$2.2$</td>
</tr>
</tbody>
</table>

### IV. CONCLUSION

In this work, the topological properties of the PESs and the molecular geometries of the XF$_3$ (X = Cl–At) fluorides have been studied with DFT and WFT based methods. From a methodological point of view, several conclusions arise. First, the PBE0 functional appears better suited to describe the whole series than the M06-2X one. Regarding the multiconfigurational WFT calculations, the ground spin-orbit-free states are found to be essentially closed-shell and single-configurational, especially in the X = Br–At cases. While spin-orbit coupling introduces some open-shell character to the AtF$_3$ wave function, the latter remains dominated by the ground spin-orbit-free spin-singlet HF configuration. It is an indisputable support for all the previous studies on these systems that made use of single-determinantal or single-reference methods.$^{1,2,4,5,9}$ Furthermore, we have demonstrated the suitability of the two-step spin-orbit configuration interaction calculations, where spin-orbit coupling is introduced \textit{a posteriori} as a perturbation of spin-orbit-free effects, for investigating
the XF$_3$ PESs, although such an approach is sometimes said to be problematic for systems containing heavy $p$-elements as At.$^{54}$

With respect to the PES explorations, new critical points have been evidenced and notably in the AtF$_3$ system, three equivalent $C_s$ TSs connecting the $C_{2v}$ minima between them. In the X = Cl–I cases, Y-shaped TSs connect two equivalent T-shaped $C_{2v}$ minima, while in the AtF$_3$ system, each Y-shaped TS connects a T-shaped $C_{2v}$ minima and the $D_{3h}$ minimum. The calculated energy difference between the $C_{2v}$ and $D_{3h}$ minima is very small, and considering the thermal and entropic contributions, no definitive conclusion can actually be given on the nature of the AtF$_3$ most stable structure. However, we firmly consider that the $D_{3h}$ minimum would be the average structure seen from experiment, if they were possible, since the Y-shaped TSs results in trifling energy barriers between the three equivalent T-shaped $C_{2v}$ minima.

Based on WFT calculations, we found in all the studied systems that scalar relativistic effects and electron correlation strongly reduce the energy differences between the T-shaped $C_{2v}$ and the $D_{3h}$ structures by more than a factor of two. While in the X = Cl–I cases, electron correlation plays a more important role, scalar relativistic effects are the main contributors to the energy differences of interest in the AtF$_3$ system. As expected in the X = Cl–I cases, no significant spin-orbit coupling effects were found, and the latter led to a weak reduction of the energy difference when X = At. Furthermore, both scalar relativistic effects and electron correlation inhibits the *pseudo* Jahn-Teller effect in the AtF$_3$ system by creating an energy barrier between the emerging $D_{3h}$ minimum and the $C_{2v}$ minima, although both effects interfere. In agreement with DFT calculations, the spin-orbit coupling effects are found negligible regarding this point.

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REFERENCES


53 See Supplementary Material Document No. for dissociation energies computed at (SO-)PBE0 levels, contributions to $\Delta E_{D_{ab}^{-C_{2v}}}$ computed at the NEVPT2 and NEVPT2+SOC levels using SA-CASSCF(2/3) references and 6S3T/3S2T state-averaging and state-interaction spaces, as well as at the CASSCF and CASSCF+SOC levels using SA-CASSCF(20/12) references and 21S20T state-averaging and state-interaction spaces.